# FURTHER WORK ON THE USE OF K, FACTORS DETERMINED BY THE CADMIUM SUBTRACTION METHOD AS A TOOL FOR A CRITICAL EVALUATION OF NUCLEAR DATA

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#### **1. INTRODUCTION**

Nuclear constants for use in reactor activation analysis and epicadmium neutron activation analysis, especially  $(n, \gamma)$  cross-sections and absolute gamma intensities, are known to show a rather large scatter in the literature [1]. This work is a continuation of that already published [2], in which shows how experimentally determined and accurate  $K_{o}$ -factors, used in a new comparator technique [5, 6], in some cases can be used to make a critical evaluation of the above mentioned constants. The method was applied to select preferred values for thermal cross-section ( $\sigma_{o}$ ), resonance integral ( $I_{o}$ ) and absolute gamma intensity.

Table	1.	Different	Reaction	and	Decay	Types	Together	With	Related	Nuclear	Parameters	and	Formulae	When	Using
							the	K₀-Me	ethod						

Reaction decay type	Reaction and decay schemes	$K_{\rm o}$ (definition)	Q <sub>o</sub> analytical	$A_{\rm sp}$ (calculation)
(I)	$1 \xrightarrow[\sigma_0, I_0]{} \stackrel{n, \nu}{\longrightarrow} 2 \xrightarrow[\rho_2]{} \stackrel{\lambda_2}{\longrightarrow}$	$\frac{M^*\theta \sigma_o \nu_2}{M \theta^* \sigma_o^* \nu^*}$	$\frac{I_{o}}{\sigma_{o}}$	$\frac{N_{\rm p,2}/t_{\rm m}}{wS_2D_2C_2}$
(II)	$1 \xrightarrow{\sigma_{o}^{m}, I_{o}^{m}} \xrightarrow{2}_{3} \xrightarrow{F_{2}, \lambda_{2}} \xrightarrow{\sigma_{o}^{g}, I_{o}^{g}} \xrightarrow{2}_{3} \xrightarrow{F_{2}, \lambda_{2}} \xrightarrow{\lambda_{2}}$	$\frac{M^*\theta \sigma_o^g \nu_3}{M \theta^* \sigma_o^* \nu^*}$	$\frac{I_o^g}{\sigma_o^g}$	$\frac{N_{\mathrm{p},3}/t_{\mathrm{m}}}{w} \left[ \frac{F_2 \sigma_{\mathrm{o}}^{\mathrm{m}}}{\sigma_{\mathrm{o}}^{\mathrm{g}}} \times \frac{(F+Q_{\mathrm{o}}^{\mathrm{m}})}{(F+Q_{\mathrm{o}}^{\mathrm{g}})} \times \frac{\lambda_3 S_2 D_2 C_2 - \lambda_2 S_3 D_3 C_3}{\lambda_3 - \lambda_2} + S_3 D_3 C_3 \right]^{-1}$
				$= \frac{N_{p,3}/t_m}{w} \left[ \frac{K_o^m}{K_o^g} \times \frac{(F+Q_o^m)}{(F+Q_o^g)} \times \frac{\lambda_3 S_2 D_2 C_2 - \lambda_2 S_3 D_3 C_3}{\lambda_3 - \lambda_2} + S_3 D_3 C_3 \right]^{-1}$
(III)	Special case: $\lambda_2 \gg \lambda_3$ and $D_2 = 0$	$\frac{M^*\theta(F_2 \sigma_o^m + \sigma_o^g) \nu_3}{M \theta^* \sigma_o^* \nu^*}$	$\frac{F_2 I_o^m + I_o^g}{F_2 \sigma_o^m + \sigma_o^g}$	$\frac{N_{\rm p,3}/t_{\rm m}}{wS_3D_3C_3}$

Target element	Atomic Weight	σ <sub>0</sub> (barn) [26]	<i>I</i> 。 (barn) [26]	Target Isotope	Abundance (%) [26]	Isotope (reaction decay type)	Half-life $T_{1/2}$	$\overline{E}_{r}$ (eV) [15]
Mg	24.30	0.063	0.038	<sup>26</sup> Mg	11.01	<sup>27</sup> Mg (I)	9.46 m	220 000
Al	26.98	0.233	0.17	<sup>27</sup> Al	100	<sup>28</sup> Al (I)	2.2405 m	8240
S	32.06	0.52	0.24	<sup>36</sup> S	0.02	<sup>37</sup> S (I)	5.06 m	
Ca	40.08	0.43	0.24	<sup>48</sup> Ca	0.187	<sup>49</sup> Ca (I)	8.72 m	
Sc	44.96	27	12	<sup>45</sup> Sc	100	<sup>46</sup> Sc (III)	83.8 d	4110
Fe	55.85	2.55	1.4	<sup>58</sup> Fe	0.28	<sup>59</sup> Fe (I)	44.50 d	325
Ga	69.72	2.9	22	<sup>71</sup> Ga	39.9	<sup>72</sup> Ga (III)	14.1 h	152
As	74.92	4.4	61	<sup>75</sup> As	100	<sup>76</sup> As (I)	26.3 h	102
Se	78.96	11.7	14	<sup>74</sup> Se	0.9	<sup>75</sup> Se (I)	119.8 d	29.5
Br	79.90	6.8	92	<sup>79</sup> Br	50.69	<sup>80m</sup> Br	4.42 h	51.4
						<sup>su</sup> Br (II)	17.6 m	
Sr	87.62	1.2	10	<sup>86</sup> Sr	9.86	<sup>87m</sup> Sr (I)	2.806 h	672

Table 2. Preferred Nuclear Data and Decay Parameters Concerning the  $(n, \gamma)$  Reaction

<b>F</b> <sub>Cd</sub>	E,	$Q_{o}(\alpha=0)$ recom. [32]	$(K_{0,Au})_{exp}$	$(K_{0,Au})_{theor}$	Evaluation value			
[5, 7, 8]	Main gamma (keV)		(rel. err., %)L		σ <sub>o</sub> (barn)	I <sub>o</sub> (barn)	Absolute gamma Intensity (%)	
1.00	170.5	0.69	$3.00 \times 10^{-6}$ (1.0)	$3.09 \times 10^{-6}$	0.038	0.025	0.86	
	843.8		$2.52 \times 10^{-4}$ (0.6)	$2.51 \times 10^{-4}$			70.0	
	1014.2		$9.79 \times 10^{-5}$ (0.6)	9.81×10 <sup>-5</sup>			27.3	
1.00	2779.0	0.73	1.77×10 <sup>-2</sup> (0.9)	$1.78 \times 10^{-2}$	0.230	0.17	100	
	3103.8	1.12	1.94×10 <sup>-6</sup> (0.9)	$1.95 \times 10^{-6}$	0.15	0.17	99.7	
	3084.2	0.45	1.00×10 <sup>-4</sup> (1.2)	$9.86 \times 10^{-5}$	1.1	0.50	92.1	
	889.3	0.44	1.25 (0.8)	1.25	27	12	100	
	1120.5		1.24 (0.4)	1.25			100	
1.00	1099.2	0.96	7.84×10 <sup>-5</sup> (0.9)	$7.86 \times 10^{-5}$	1.33	1.1	56.5	
	1291.6		5.99×10 <sup>-5</sup> (0.9)	$6.01 \times 10^{-5}$			43.2	
1.00	629.9	6.63	1.48×10 <sup>-2</sup> (0.8)	1.46×10 <sup>-2</sup>	4.65	30.8	26.3	
	834.1		$5.27 \times 10^{-2} (0.5)$	$5.31 \times 10^{-2}$			95.54	
	894.2		5.49×10 <sup>-3</sup> (0.9)	$5.47 \times 10^{-3}$			9.847	
	1050.7		$3.85 \times 10^{-3}$ (1.1)	$3.84 \times 10^{-3}$			6.921	
	2201.6		$1.44 \times 10^{-2}$ (1.0)	$1.45 \times 10^{-2}$			26.07	
	2490.9		$4.20 \times 10^{-3}$ (1.5)	$4.15 \times 10^{-3}$			7.476	
	2501.6 $(E_{eff.})$		$1.15 \times 10^{-2}$ (1.3)	$1.14 \times 10^{-2}$			20.54	
	2507.7 ( $E_{\rm eff.}$ )		7.30×10 <sup>-3</sup> (1.0)	$7.25 \times 10^{-3}$			13.06	
1.00	559.2 $(E_{eff.})$	14.0	4.99×10 <sup>-2</sup> (1.0)	$5.03 \times 10^{-2}$	4.3	60	41.99	
	657.0		$6.65 \times 10^{-3} (0.8)$	$6.82 \times 10^{-3}$			5.69	
	1215.2 $(E_{\rm eff.})$		$5.25 \times 10^{-3}$ (1.2)	$5.21 \times 10^{-3}$			4.35	
0.94	121.1	10.9	1.98×10 <sup>-3</sup> (0.5)	$1.99 \times 10^{-3}$	51.8	565	16.15	
	136.0		$6.88 \times 10^{-3} (0.6)$	$6.90 \times 10^{-3}$			56.02	
	264.7		$7.23 \times 10^{-3}$ (0.9)	$7.22 \times 10^{-3}$			58.60	
	279.5		$3.06 \times 10^{-3} (0.9)$	$3.07 \times 10^{-3}$			24.9	
	400.6		1.43×10 <sup>-3</sup> (0.8)	$1.45 \times 10^{-3}$			11.78	
1.00								
1.00	616.9	11.4	$6.68 \times 10^{-3} (0.9)$	$7.01 \times 10^{-3}$	0.4	96	6.3	
	666.3		1.18×10 <sup>-3</sup> (0.9)	$1.17 \times 10^{-3}$			1.05	
1.00	388.5	4.11	1.49×10 <sup>-3</sup> (0.7)	$1.49 \times 10^{-3}$	0.369	3.16	82.5	

### 2. EXPERIMENTAL DETAILS AND RESULTS

The experimental method has been described previously [2-4]; in the present work all  $K_0$  factors were experimentally determined and calculated versus the 411.794 keV  $\gamma$ -line of <sup>198</sup>Au [5]. The nuclear data of interest, which are required for the calculations, are listed in Table 1 of Reference [2]. Counting was performed using a 104 cm<sup>3</sup> and 40 cm<sup>3</sup> single open-ended coaxial Ge(Li) detector with a source-to-detector separation of 15 cm. Appropriate formulae for specific count rate calculations (and  $K_{0}$ and  $Q_{o}$  definitions) for several complex decay schemes (mother-daughter decay, etc.) are given in Table 1. The listed  $K_o$ -factors are those recommended when obtained according to the experimental method that has been described previously [2] and when the standard deviation on the mean was less than 2%. It should be realized that the average usually results from five measurements on two irradiation channels. Results of  $K_{0}$ determinations are shown in Table 2 for 12 isotopes, covering 29  $\gamma$ -lines. In this evaluation several compilations were systematically surveyed [9-31]and other recent publications were consulted. The finally adopted values are based on the present evaluation result, and *t*-test can be performed for the comparison of the theoretical  $(X)_{\text{theor}}$  and mean experimental value  $(\bar{X})_{exp}$ . If  $S_{exp}$  is the standard deviation on the mean for an average value  $(X)_{exp}$ , then

$$t = \frac{(\bar{X})_{\exp} - X_{\text{theor}}}{S_{\exp}}$$

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